Symmetric α -Stable Random Variables, Fractional Calculus, and the Fourier-Mellin Triangle

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- This talk is about the interaction between:
 - the probability density function of symmetric α-stable (SαS) random variables;
 - 2 the Fourier transform;
 - and the Mellin transform.

This Talk in a Nutshell

- This interaction is mediated by techniques from fractional calculus.
- Many of these ideas have been developed by Di Paolo et al in
 - Cottone and Di Paolo (2009), On the use of fractional calculus for the probabilistic characterization of random variables.
 - ② Di Paolo and Pinnola (2012), Riesz fractional integrals and complex fractional moments for the probabilistic characterization of random variables.

This Talk in a Nutshell

• The following diagram—the Fourier-Mellin triangle—plays a key role:



Why α -Stable Random Variables?

- Distributions with heavy tails are not well-modeled as Gaussian.
- α -stable distributions are a relatively tractable alternative.
 - $\mathbb{P}(X > \lambda) \sim C\lambda^{-\alpha}$
- Applications:
 - Interference and noise modeling (e.g., in wireless radio communications).
 - Asset returns in finance.

Symmetric α -Stable Random Variables

- An important sub-class are the symmetric α-stable (SαS) random variables.
- One way $S\alpha S$ random variables arise is the LePage series

$$X = \sum_{i=1}^{\infty} r_i^{-\alpha} g_i,$$

where

- **(** $\{r_i\}$ are the arrival times of the Poisson process with rate 1;
- {g_i} are symmetric random variables X ^d/₌ −X, independent of {r_i};
 𝔅[g_i^α] < ∞.
- In applications this type of sum is called *shot noise*.

Symmetric α -Stable Random Variables

• A common definition of $S\alpha S$ random variables is via

$$AX_1 + BX_2 \stackrel{d}{=} CX, \quad X \stackrel{d}{=} -X,$$

where

X₁, X₂ are independent copies of X;
 A^α + B^α = C^α, for some α ∈ (0,2].

Symmetric α -Stable Random Variables

- $S\alpha S$ random variables are also infinitely divisible.
- For $n \ge 2$, there is a $C_n > 0$ such that

$$X_1 + \cdots X_n \stackrel{d}{=} C_n X, \quad X \stackrel{d}{=} -X,$$

where

Examples of Symmetric α -Stable Random Variables

1 Gaussian distribution
$$(X \sim S_2(\sigma, 0, 0))$$
:

$$p_X(x) = \frac{1}{\sqrt{4\pi\sigma^2}} e^{-\frac{x^2}{4\sigma^2}}$$
$$\mathbb{E}[e^{itX}] = e^{-\sigma^2 t^2}$$

2 Cauchy distribution
$$(X \sim S_1(\sigma, 0, 0))$$
:

$$p_X(x) = \frac{1}{\pi\sigma} \left[\frac{\sigma^2}{x^2 + \sigma^2} \right]$$
$$\mathbb{E}[e^{itX}] = e^{-\sigma|t|}$$

In general, $S\alpha S$ random variables do not have closed form densities.

The Characteristic Function of $S\alpha S$ Random Variables

• While the density of $S\alpha S$ random variables is difficult to work with, the characteristic function

$$\mathbb{E}[e^{itX}] = \int_{-\infty}^{\infty} p_X(x) e^{itX} dx$$

is known in closed form.

• In particular, let $X \sim S_{lpha}(\sigma, 0, 0)$. Then,

$$\phi_X(t) = \mathbb{E}[e^{itX}] = e^{-\sigma^{\alpha}|t|^{\alpha}}$$

• This result is very useful.

Basic Properties of $S\alpha S$ Random Variables

1 Let X_1, X_2 be independent with $X_i \sim S_{\alpha}(\sigma_i, 0, 0)$. Then,

$$X_1 + X_2 \sim S_{\alpha} \left((\sigma_1^{\alpha} + \sigma_2^{\alpha})^{1/\alpha}, 0, 0 \right).$$

2 Let $X \sim S_{\alpha}(\sigma, 0, 0)$ and $a \in \mathbb{R}$. Then,

$$X + a \sim S_{\alpha}(\sigma, 0, a).$$

• Let $X \sim S_{\alpha}(\sigma, 0, 0)$ and $a \in \mathbb{R} \setminus \{0\}$. Then,

 $aX \sim S_{\alpha}(|a|\sigma, 0, 0).$

Basic Properties of $S\alpha S$ Random Variables

• Let $X \sim S_{\alpha}(\sigma, 0, 0)$. Then,

$$\mathbb{P}(X > \lambda) \sim \sigma^{lpha} rac{\mathcal{C}_{lpha}}{2} \lambda^{-lpha}.$$

2 Let $X \sim S_{\alpha}(\sigma, 0, 0)$ and 0 . Then,

$$\mathbb{E}[|X|^{p}] = \frac{2^{p+1}\Gamma\left(\frac{p+1}{2}\right)\Gamma\left(-\frac{p}{\alpha}\right)}{\alpha\sqrt{\pi}\Gamma\left(-\frac{p}{2}\right)}\sigma^{p}.$$

- Instead of taking the Fourier transform (i.e., the characteristic function), we can consider the Mellin transform.
- For α -stable distributions, this was first done by Zolotarev in 1957.
- Since SαS densities are absolutely continuous functions, the Mellin transform is

$$\mathcal{M}[p_X(x)](\gamma) = \int_0^\infty p_X(x) x^{\gamma-1} dx, \ \gamma \in \mathbb{C}.$$

(More generally, the Mellin-Stieltjes transform is required.)

- The Mellin transform can be related to the Fourier transform through a change of variables.
- Consider the operator $T : f(x) \rightarrow f(e^x), f \in L^1$; i.e.,

$$\int_{-\infty}^{\infty} |f(x)| dx < \infty.$$

Define the T-norm

$$\|f\|_{\mathcal{M}_{c}} = \int_{0}^{\infty} |f(x)| x^{c-1} | dx, \qquad (1)$$

where c is chosen to ensure convergence for the class of functions f we are interested in.

• L¹ functions with finite T-norm

$$\int_0^\infty |f(x)| x^{c-1} | dx < \infty \tag{2}$$

form a function space.

• On this space, the Fourier transform

•

$$\mathcal{F}(Tf)(t) = \int_{-\infty}^{\infty} f(e^x) e^{itx} dx.$$

is well defined. (More on this in Butzer and Jansche (1997)).

Now consider

$$\mathcal{F}[Tf]\left(\frac{\eta-c}{i}\right) = \int_{-\infty}^{\infty} f(e^x) e^{(\eta-c)x} dx, \ c \ge 0.$$

• Let $y = e^x$, which gives

$$\mathcal{F}(Tf)\left(\frac{\eta-c}{i}\right) = \int_0^\infty f(y)e^{(\eta-c)\log y}\frac{1}{y}dy$$
$$= \int_0^\infty f(y)y^{-c}y^{\eta}\frac{1}{y}dy$$
$$= \int_0^\infty f^*(y)y^{\eta-1}dy, \ f^*(y) = f(y)y^{-c},$$

which is the Mellin transform of $f^*(y)$.

• This means that for functions with finite T-norm

$$\|f\|_{\mathcal{M}_{c}} = \int_{0}^{\infty} |f(x)| x^{c-1} dx,$$
(3)

we can use theorems for the Fourier transform of L^1 functions.

• E.g., the Fourier inversion theorem (more on this later).

- The Mellin transform has a nice property.
- Consider the product of two random variables Z = XY. Then,

$$\mathcal{M}[f_Z](s) = \mathcal{M}[f_X](s)\mathcal{M}[f_Y](s). \tag{4}$$

• This means that the Mellin transform plays a similar role for products of random variables as the Fourier transform plays for sums.

- The Mellin transform has an intimate link to fractional calculus.
- To see this, we first overview some basic ideas in fractional calculus.

• The starting point for fractional calculus is to generalize derivatives

$$rac{d^n}{dx^n}f(x), \ n\in\mathbb{N}$$

to the case where $n \in \mathbb{R}$; e.g.,

$$\frac{d^{\frac{1}{2}}}{dx^{\frac{1}{2}}}f(x).$$

- To see how this might work, consider $f(x) = x^{p}$.
- We have,

$$\frac{d^n}{dx^n}x^p = p(p-1)\cdots(p-n+1)x^{p-n}$$
$$= \frac{p!}{(p-n)!}x^{p-n} = \frac{\Gamma(p+1)}{\Gamma(p-n+1)}x^{p-n}.$$

• Using the properties of the Gamma function

$$\Gamma(x)=\int_0^\infty e^{-t}t^{x-1}dt,$$

it is possible to analytically continue to yield

$$\frac{d^q}{dx^q}x^p = \frac{\Gamma(p+1)}{\Gamma(p-q+1)}x^{p-q}$$

with $q \in \mathbb{R}$ (being careful with q = -1, -2, ...).

• An important way of generalizing is via the Riemann-Liouville fractional integrals

$$(I_{\pm}^{\gamma}f)(x) = rac{1}{\Gamma(\gamma)}\int_{0}^{\infty}\zeta^{\gamma-1}f(x\mp\zeta)d\zeta.$$

• The Riemann-Liouville fractional derivatives are then

$$(D_{\pm}^{\gamma}f)(x) = \frac{(\pm 1)^n}{\Gamma(n-\gamma)} \frac{d^n}{dx^n} \int_0^{\infty} \zeta^{n-\gamma-1} f(x \mp \zeta) d\zeta,$$

where $\gamma \in \mathbb{C}$ and $n = [\rho] + 1$, where $\rho = \operatorname{Re}(\gamma)$.

 This agrees with gamma function-based definition for the monomial example developed in the previous slide.

• The Riemann-Liouville fractional integral satisfies the semigroup property:

$$I^{\alpha}_{+}I^{\beta}_{+}\psi = I^{\alpha+\beta}_{+}\psi, \quad I^{\alpha}_{-}I^{\beta}_{-}\psi = I^{\alpha+\beta}_{-}\psi, \tag{5}$$

where $\alpha, \beta > 0$.

The Link Between Fractional Calculus and the Mellin Transform

- A key observation is the link between the Riemann-Liouville fractional integral and the Mellin transform.
- Recall the Mellin transform is

$$\mathcal{M}[f(x)](\gamma) = \int_0^\infty p_X(x) x^{\gamma-1} dx, \ \gamma \in \mathbb{C}.$$

• The Riemann-Liouville fractional integral is

$$(I_{\pm}^{\gamma}f)(x) = rac{1}{\Gamma(\gamma)}\int_{0}^{\infty}\zeta^{\gamma-1}f(x\mp\zeta)d\zeta.$$

That is,

$$\mathcal{M}[f(x \mp \zeta)](\gamma) = \Gamma(\gamma)(I_{\pm}^{\gamma}f)(x).$$

A Key Identity

• For standard integrals, we have

$$\mathcal{F}\left[\int_{-\infty}^{x} f(\tau) d\tau\right](t) = rac{\mathcal{F}[f](t)}{-it}$$

• This generalizes:

$$\mathcal{F}[I^{\gamma}_{+}f](t) = (-it)^{-\gamma}\mathcal{F}[f](t).$$

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A Key Identity

 Take 0 < ρ < 1. We need to take the Fourier transform of the fractional integral; i.e.,

$$\mathcal{F}[l_+^{\gamma}f](t) = \int_{-\infty}^{\infty} e^{itx} \frac{1}{\Gamma(\gamma)} \int_0^{\infty} \zeta^{\gamma-1} f(x-\zeta) d\zeta dx.$$

Some basic manipulations yield

$$\mathcal{F}[l_+^\gamma f](t) = rac{\mathcal{F}[f](t)}{ \mathsf{\Gamma}(\gamma)} \int_0^\infty e^{it\zeta} \zeta^{\gamma-1} d\zeta.$$

• A useful identity tells us that

$$\int_0^\infty e^{it\zeta}\zeta^{\gamma-1} = \Gamma(\gamma)(-it)^{-\gamma},$$

where taking the principle value we understand that

$$(-it)^{-\gamma} = \exp\left(-\gamma \log|t| + \frac{\gamma \pi i}{2} \operatorname{sgn}(t)\right).$$

• Leading us to

$$\mathcal{F}[I^{\gamma}_{+}f](t) = (-it)^{-\gamma} \mathcal{F}[f](t).$$

• What does it mean?

Image: A math a math

Recovering the Mellin Transform

• The observation that

$$\mathcal{F}[I^{\gamma}_{+}f](t) = (-it)^{-\gamma}\mathcal{F}[f](t).$$

means that

$$egin{aligned} &(I^{\gamma}_+f)(t) = \mathcal{F}^{-1}\left[(-it)^{-\gamma}\mathcal{F}[f](t)
ight] \ &= rac{1}{2\pi}\int_{-\infty}^{\infty}e^{-itx}(-it)^{-\gamma}\mathcal{F}[f](t)dt. \end{aligned}$$

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Recovering the Mellin Transform

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• As such,

$$(I^{\gamma}_{+}f)(0)=rac{1}{2\pi}\int_{-\infty}^{\infty}(-it)^{-\gamma}\mathcal{F}[f](t)dt.$$

Now recall that

$$\Gamma(\gamma)(I_+^{\gamma}f)(x) = \int_0^{\infty} \zeta^{\gamma-1}f(x-\zeta)d\zeta.$$

Recovering the Mellin Transform

• Putting it all together:

$$\int_0^{\infty} \zeta^{\gamma-1} f(-\zeta) d\zeta = \frac{\Gamma(\gamma)}{2\pi} \int_{-\infty}^{\infty} (-it)^{-\gamma} \mathcal{F}[f](t) dt.$$

• When f is even, we then have

$$\mathcal{M}[f](\gamma) = \int_0^\infty \zeta^{\gamma-1} f(\zeta) d\zeta$$

= $\frac{\Gamma(\gamma)}{2\pi} \int_{-\infty}^\infty (-it)^{-\gamma} \mathcal{F}[f](t) dt.$

The Mellin Transform and the $S\alpha S$ Characteristic Function

• We now choose f to be a $S\alpha S$ density p_X , with

$$\phi_X(t) = e^{-\sigma^lpha |t|^lpha}$$

This means that

$$\mathcal{M}[p_X](\gamma) = \int_0^\infty \zeta^{\gamma-1} p_X(\zeta) d\zeta$$
$$= \frac{\Gamma(\gamma)}{2\pi} \int_{-\infty}^\infty (-it)^{-\gamma} \phi_X(t) dt$$

• Using the fact that $\phi_X(t)$ is real and $\phi_X(t)^* = \phi_X(-t)$, it follows that

$$\mathcal{M}[p_X](\gamma) = \frac{\Gamma(\gamma)\cos\left(\frac{\gamma\pi}{2}\right)}{\pi} \int_0^\infty t^{-\gamma}\phi_X(t)dt.$$

The Fourier-Mellin Triangle

• Going a step further, we can identify

$$\int_{0}^{\infty} t^{-\gamma} \phi_{X}(t) dt = \mathcal{M}[\phi_{X}](1-\gamma)$$
$$\Rightarrow \mathcal{M}[p_{X}](\gamma) = \frac{\Gamma(\gamma) \cos\left(\frac{\gamma \pi}{2}\right)}{\pi} \mathcal{M}[\phi_{X}](1-\gamma).$$

The Fourier-Mellin Triangle

• This all can now be summarized by the Fourier-Mellin triangle:



where

$$\mathcal{G}[\phi_X](\gamma) = rac{\mathsf{F}(\gamma)\cos\left(rac{\gamma\pi}{2}
ight)}{\pi}\mathcal{M}[\phi_X](1-\gamma).$$

The Mellin Transform of $S\alpha S$ Densities

- We can use the Fourier-Mellin triangle to evaluate the Mellin transform of $S\alpha S$ densities.
- In particular, we have

$$M_X(\gamma) = \frac{\Gamma(\gamma)\cos\left(\frac{\gamma\pi}{2}\right)}{\pi} \int_0^\infty e^{-\sigma^\alpha t^\alpha} t^{-\gamma} dt$$
$$= \frac{\sigma^{\gamma-1}\Gamma(\gamma)\Gamma\left(\frac{1-\gamma}{\alpha}\right)}{\pi\alpha}\cos\left(\frac{\gamma\pi}{2}\right).$$

- Note that this method generalizes to any symmetric distribution, and can also be further generalized to asymmetric distributions.
- See di Paolo and Pinnola (2012) for more details.

Recovering the Density: The Inverse Mellin Transform

- The Fourier-Mellin triangle provides a convenient way to obtain the Mellin transform.
- We can use the Mellin transform of $S\alpha S$ densities to recover the density.
- This approach has been developed in Cottone and Di Paolo (2009) and Di Paolo and Pinnola (2012).
- The relevant tool is the inverse Mellin transform:

$$p_X(x) = rac{1}{2\pi i} \int_{
ho-i\infty}^{
ho+i\infty} M_X(\gamma) |x|^{-\gamma} d\gamma, \ x
eq 0.$$

• Condition:

1 γ must lie in the fundamental strip.

The Fundamental Strip

- The fundamental strip is the set of $\rho = \operatorname{Re}(\gamma)$ for which the Mellin integral converges.
- To see when this occurs for the Mellin transform

$$M_X(\gamma) = \int_0^\infty p_X(x) x^{\gamma-1} dx,$$

we can use the Fourier-Mellin triangle; i.e.,

$$M_X(\gamma) = \frac{\Gamma(\gamma) \cos\left(\frac{\gamma \pi}{2}\right)}{\pi} \mathcal{M}[\phi_X](1-\gamma).$$

The Fundamental Strip

• In particular, observe that

$$|\mathcal{M}[\phi_X](1-\gamma)| \leq \int_0^1 t^{-
ho} dt + \int_1^\infty |\phi_X(t)| dt.$$

• Since

$$\int_0^\infty e^{-\sigma^\alpha t^\alpha} dt = \frac{\Gamma\left(\frac{1}{\alpha}\right)}{\alpha \sigma^\alpha},$$

it follows that the Mellin transform converges for 0 < ρ < 1.

• The integral density representation

$$egin{aligned} p_X(x) &= rac{1}{2\pi i} \int_{
ho-i\infty}^{
ho+i\infty} M_X(\gamma) |x|^{-\gamma} d\gamma, \; x
eq 0 \ &= rac{1}{2\pi} \int_{-\infty}^{\infty} M_X(
ho+i\eta) |x|^{-
ho-i\eta} d\eta, \; x
eq 0, \end{aligned}$$

lends itself to approximation.

• In particular, we can use the trapezoidal approximation

$$p_X(x) \approx \frac{\Delta \eta}{2\pi} \sum_{k=-m}^m M_X(\gamma_k) |x|^{-\gamma_k},$$

where $\gamma_k = \rho + ik\Delta\eta$.

- Di Paolo and Pinnola (2012) investigated the trapezoidal approximation.
- For the symmetric Cauchy density ($\alpha = 1$), they found that for $\sigma = 0.6$, choosing $\Delta \eta = 0.4$, $\rho = 0.5$ leads to



• For the symmetric Gaussian density ($\alpha = 2$), Cottone and di Paolo (2009) found that for $\sigma^2 = 1$, choosing $\Delta \eta = 0.4$, $\rho = 0.4$, leads to



- However, there is a problem when other approximation parameters are chosen.
- E.g., symmetric Cauchy with $\sigma=$ 0.6, choosing $\Delta\eta=$ 0.4, $\rho=$ 0.2 (vs $\rho=$ 0.5) leads to



Bounding the Approximation Error

- To overcome this problem, we need error bounds.
- There are two sources of error:
 - Truncation error.
 - 2 Discretization error.



Bounding the Truncation Error

• For the truncation error, we need to bound the integral

$$\begin{split} |E_{\mathcal{T},\mathcal{R}}| &= \left|\frac{1}{2\pi}\int_{m\Delta\eta}^{\infty}M_{X}(\rho+i\eta)|x|^{-\rho-i\eta}d\eta\right| \\ &\leq \frac{1}{2\pi}\int_{m\Delta\eta}^{\infty}\left|\frac{\Gamma(\rho+i\eta)\Gamma\left(\frac{1-\rho-i\eta}{\alpha}\right)}{\pi\alpha}\cos\left(\frac{(\rho+i\eta)\pi}{2}\right)|x|^{-\rho-i\eta}\right|d\eta \\ &\leq \frac{1}{2\pi^{2}\alpha}\int_{m\Delta\eta}^{\infty}\left|\Gamma(\rho+i\eta)\Gamma\left(\frac{1-\rho-i\eta}{\alpha}\right)\cosh\left(\frac{\pi\eta}{2}\right)\right||x|^{-\rho}d\eta. \end{split}$$

• In the case $\alpha = 1$, we can use

$$\Gamma(z)\Gamma(1-z) = \frac{\pi}{\sin(\pi z)}.$$

Bounding the Truncation Error

• This leads to

$$\begin{split} E_{\mathcal{T}} &| \leq \frac{|x|^{-\rho}}{2\pi} \left(\int_{m\Delta\eta}^{\infty} \left| \frac{\cosh\left(\frac{\pi\eta}{2}\right)}{\sin(\pi(\rho+i\eta))} \right| d\eta \\ &+ \int_{-\infty}^{-m\Delta\eta} \left| \frac{\cosh\left(\frac{\pi\eta}{2}\right)}{\sin(\pi(\rho+i\eta))} \right| d\eta \right) \\ &\leq \frac{|x|^{-\rho}\sqrt{2}}{2\pi} \left(\int_{m\Delta\eta}^{\infty} \frac{\cosh\left(\frac{\pi\eta}{2}\right)}{\cosh(\pi\eta)} d\eta + \int_{-\infty}^{-m\Delta\eta} \frac{\cosh\left(\frac{\pi\eta}{2}\right)}{\cosh(\pi\eta)} d\eta \right) \end{split}$$

• Observe that for |x| > 1, the truncation error bound improves for larger ρ .

• We now turn to the discretization error.



- The discretization error of the trapezoidal rule can be obtained using the residue theorem.
- Define

$$I_{\Delta\eta}(x) = \frac{\Delta\eta}{2\pi} \sum_{k=-\infty}^{\infty} M_X(\gamma_k) |x|^{-\gamma_k},$$

where $\gamma_k = \rho + i\Delta\eta$.

• The discretization error is then

$$DE = |p_X(x) - I_{\Delta\eta}(x)|.$$

• Under certain regularity conditions, the trapezoidal rule has

$$DE \leq rac{2M}{e^{2\pi c/h}-1},$$

where c is a bound on the analytic region of the function being integrated.

• In our case, the function is

$$M_X(\gamma)|x|^{-\gamma}.$$
 (6)

As such,

$$DE \leq rac{2M}{e^{2\pi(1-
ho)/\Delta\eta}-1},$$

where

$$\int_{-\infty}^{\infty} |M_X(\rho + iu - r)|x|^{-\rho - iu + r}| du \leq M,$$

for all $\rho - 1 < r < 0$.

• A key point is that the discretization error decays as $O(e^{-2\pi(1-\rho)/\Delta\eta})$.

• In both the truncation error and the discretization error, the term

 $|x|^{-\rho}$

appears on the numerator.

- This means that the approximation *improves* for large |x|.
- That is, this method can be useful for approximating the tails of SαS densities.

More On Fractional Calculus

 Earlier in the talk the Riemann-Liouville fractional derivative was introduced as

$$(D_{\pm}^{\gamma}f)(x) = \frac{(\pm 1)^n}{\Gamma(n-\gamma)} \frac{d^n}{dx^n} \int_0^{\infty} \zeta^{n-\gamma-1} f(x \mp \zeta) d\zeta,$$

where $\gamma \in \mathbb{C}$ and $n = [\rho] + 1$, where $\rho = \operatorname{Re}(\gamma)$.

Another type of fractional derivative is due to Riesz, given by

$$(D^{\gamma}f)(x) = -\frac{1}{2\cos(\gamma\pi/2)}((D^{\gamma}_{+}f)(x) + (D^{\gamma}_{-}f)(x)).$$

More on Fractional Calculus

- The Riesz fractional derivative has a strong link to fractional moments.
- Cottone and Di Paolo (2009) have shown that

$$(D^{\gamma}\phi_X)(0) = -\mathbb{E}[|X|^{\gamma}], \operatorname{Re}(\gamma) > 0.$$

• This can be viewed analogously to the usual result

$$\mathbb{E}[X^n] = i^{-n} \phi_X^{(n)}(0). \tag{7}$$

More on Fractional Calculus

- To prove it, take the Fourier transform of the Riemann-Liouville fractional derivative.
- This yields

$$\mathcal{F}[(D_{\pm}^{\gamma}\phi_X)](x) = (\mp ix)^{\gamma}\mathcal{F}[\phi_X](x),$$

analogous to the integer derivative case.

• Taking the inverse Fourier transform and setting t to zero yields

$$(D_{\pm}^{\gamma}\phi_X)(0)=\mathbb{E}[(\mp iX)^{\gamma}].$$

• The result

$$(D^{\gamma}\phi_X)(0) = -\mathbb{E}[|X|^{\gamma}], \operatorname{Re}(\gamma) > 0.$$

then follows by straightforward manipulations.

Fractional Moments

• A standard result is that the fractional moments for $S\alpha S$ random variables are given by

$$\mathbb{E}[|X|^{p}] = \frac{2^{p+1} \Gamma\left(\frac{p+1}{2}\right) \Gamma\left(-\frac{p}{\alpha}\right)}{\alpha \sqrt{\pi} \Gamma\left(-\frac{p}{2}\right)} \sigma^{p}.$$

• Sometimes it is also useful to compute moments of the form:

$$\mathbb{E}[|X-\mu|^p].$$

• Why? Consider the α -stable noise channel

$$Y = X + N,$$

where $N \sim S_{\alpha}(\sigma, 0, 0)$.

- This is a useful model for interference in large scale wireless communication networks.
- A key step in deriving an upper bound on the capacity

 $\max_{r_X} I(X;Y)$

is to compute moments of the form $\mathbb{E}[|X - \mu|^p]$.

Fractional Moments

- The problem of finding fractional moments E[|X − μ|^p] has been studied by Matsui and Pawlas (2014) in the case α > 1.
- Their approach relied on the use of the Marchaud fractional derivative

$$\frac{d^{\gamma}}{dt^{\gamma}}f(t)=\frac{\lambda}{\Gamma(1-\lambda)}\int_{-\infty}^{t}\frac{f^{(k)}(t)-f^{(k)}(u)}{(t-u)^{1+\lambda}}du,\ t\in\mathbb{R},$$

where $\gamma = k + \lambda$, with $k \in \mathbb{N}$ and $0 < \lambda < 1$.

Fractional Moments

• Let
$$m_{\mu,1+\lambda} = \mathbb{E}[|X - \mu|^{1+\lambda}].$$

• Matsui and Pawlas (2014) showed that for $S\alpha S$ random variables with $1 < 1 + \lambda \leq 2$

$$m_{\mu,1+\lambda} = \frac{\lambda \sigma^{1+\lambda}}{\sin\left(\frac{\lambda \pi}{2}\right) \Gamma(1-\lambda)} \left[\frac{\mu}{\sigma} \int_0^\infty u^{-(1+\lambda)} e^{-u^\alpha} \sin\left(\frac{\mu u}{\sigma}\right) du \right. \\ \left. + \alpha \int_0^\infty u^{\alpha-\lambda-2} e^{-u^\alpha} \cos\left(\frac{\mu u}{\sigma}\right) du \right].$$

• We want to extend this result to the case 0 $< 1 + \lambda < 1.$

• To do this we use the identity of Cottone and Di Paolo (2009)

$$(D^{\gamma}\phi_X)(0) = -\mathbb{E}[|X|^{\gamma}], \operatorname{Re}(\gamma) > 0$$

and the definition of the Riesz fractional derivative

$$(D^{\gamma}f)(x) = -\frac{1}{2\cos(\gamma\pi/2)}((D^{\gamma}_{+}f)(x) + (D^{\gamma}_{-}f)(x)).$$

 This means we only need to compute the individual Riemann-Liouville fractional derivatives

$$(D_{\pm}^{\gamma}f)(x) = \frac{(\pm 1)^n}{\Gamma(n-\gamma)} \frac{d^n}{dx^n} \int_0^{\infty} \zeta^{n-\gamma-1} f(x \mp \zeta) d\zeta,$$

Fractional Moments

• In our case, we need to compute

$$(D^{\gamma}_{\pm}f)(x) = rac{\pm 1}{\Gamma(1-\gamma)} rac{d}{dx} \int_{0}^{\infty} \zeta^{-\gamma} e^{i\mu t - \sigma^{lpha} |x \mp \zeta|^{lpha}} d\zeta,$$

• This yields

$$\mathbb{E}[|X + \mu|^{p}] = \frac{\sigma^{p}}{\Gamma(1 - p)\cos(p\pi/2)} \left[\mu \int_{0}^{\infty} u^{-p} e^{-u^{\alpha}} \sin(\mu u/\sigma) du + \alpha \int_{0}^{\infty} u^{\alpha - p - 1} e^{-u^{\alpha}} \cos(\mu u/\sigma) du \right].$$

A Summary So Far

- So far for univariate $S\alpha S$ random variables, we have looked at:
 - Interstation of the second second
 - Obtained the Fourier-Mellin triangle



3 Using the integral representation to approximate $S\alpha S$ densities. 3 Derived fractional moments $\mathbb{E}[|X - \mu|^p]$.

Extensions to Multivariate $S\alpha S$

- It is also possible to extend a number of the results to the multivariate setting.
- This is achieved via the multivariable Mellin transform and multivariable fractional calculus.

Multivariate $S\alpha S$ Random Vectors

- The definition of $S\alpha S$ random variables can be generalized.
- I.e., a random vector $X = (X_1, ..., X_n)$ is $S \alpha S$ if for any A, B > 0, there is a C > 0 and vector $\in \mathbb{R}^n$ such that

$$A\mathbf{X}^{(1)} + B\mathbf{X}^{(2)} \stackrel{d}{=} C\mathbf{X} + \mathbf{d},$$
(8)

and $\mathbf{X} \stackrel{d}{=} -\mathbf{X}$.



The Characteristic Function

- For $S\alpha S$ random vectors, we can write the characteristic function in two ways.
- In the bivariate case, we have

2

1

$$\phi(t_1,t_2)=\int_{\mathbb{R}^2}e^{i\mathbf{t}\cdot\mathbf{x}}p(x_1,x_2)dx_1dx_2.$$

$$\phi(t_1, t_2) = \exp\left\{-\int_{\mathbb{S}^{d-1}}\left|\sum_k t_k s_k\right|^{\alpha} d\Gamma(s_1, s_2)\right\}.$$

The Multivariate Mellin Transform

• By exploiting the link with the Fourier transform, we can generalize the Mellin transform to \mathbb{R}^d as

$$\mathcal{M}[f](\mathbf{s}) = \int_0^\infty \cdots \int_0^\infty f(\zeta) \prod_{i=1}^d \zeta_i^{\mathbf{s}_i - 1} d\zeta.$$
(9)

The Multivariate Riemann-Liouville Fractional Integral

 Similarly, we can define the multivariable Riemann-Liouville fractional integral as

$$(I_{\pm}^{\gamma}f)(\mathbf{x}) = \frac{1}{\Gamma(\alpha)} \int_0^{\infty} \cdots \int_0^{\infty} f(x \mp \zeta) \prod_{i=1}^d \zeta_i^{\gamma_i - 1} d\zeta, \qquad (10)$$

which is again closely linked to the Mellin transform as

$$\Gamma(\alpha)(I_{\pm}^{\gamma}f)(\mathbf{x}) = \int_0^{\infty} \cdots \int_0^{\infty} f(x \mp \zeta) \prod_{i=1}^d \zeta_i^{\gamma_i - 1} d\zeta.$$
(11)

The Multivariate Riemann-Liouville Fractional Integral

• We can also consider the Fourier transform of $(I_{\pm}^{\gamma}f)(\mathbf{x})$, which is

$$\mathcal{F}[(I_{\pm}^{\gamma}f)(\mathbf{x})](\mathbf{t}) = \frac{\mathcal{F}[f](t)}{\prod_{j=1}^{d} (\mp it)^{\gamma_{i}}}.$$
(12)

This yields

$$(I_{\pm}^{\gamma}f)(0) = \mathbb{E}[\prod_{j=1}^{d} (\mp iX_j)^{-\gamma_j}].$$
 (13)

The Multivariate Riemann-Liouville Fractional Integral

• By applying these results to the characteristic function and taking the inverse Mellin transform, we obtain

$$\phi(\pm\zeta) = \frac{1}{2\pi i} \int_{\rho+i\mathbb{R}^d} \Gamma(\gamma) \mathbb{E}[\prod_{j=1}^d (\mp iX_j)^{-\gamma_j}] \zeta^{-\gamma} d\gamma, \ \zeta \succ 0.$$
(14)

- This yields a third way of representing the characteristic function.
- This is important because it implies that the characteristic function is completely described by the fractional moment surface

$$\mathbb{E}\left[\prod_{j=1}^{d} (\mp i X_j)^{-\gamma_j}\right]$$
(15)

A Question

- In the multivariable case, there is the issue of dependence.
- This can be characterized in two ways:
 - In the density $p(x_1, x_2)$, for instance, using a copula:

$$p(x_1, x_2) = p_1(x_1)p_2(x_2)c(F_1(x_1), F_2(x_2)).$$

2 In the spectral measure Γ .

How can we relate the spectral measure $\Gamma(s_1, s_2)$ to the dependence structure in the density $p(x_1, x_2)$?

• The representation

$$\begin{split} \phi(t_1, t_2) &= \frac{1}{(2\pi i)^2} \int_{\rho + i\mathbb{R}^2} \Gamma(\gamma_1) \Gamma(\gamma_2) \\ &\times \mathbb{E}[(-iX_1^{-\gamma_1})(-iX_2^{-\gamma_2})] t_1^{-\gamma_1} t_2^{-\gamma_2} dt_1 dt_2. \end{split}$$

suggests that this might be possible by studying the surface

$$\mathbb{E}[(-iX_1^{-\gamma_1})(-iX_2^{-\gamma_2})],$$

which completely characterizes the random vector.

• This remains on-going work.

Conclusions

- We looked at univariate $S\alpha S$ random variables, where we:
 - Observed the link between the Mellin transform and fractional calculus.
 - Obtained the Fourier-Mellin triangle



- **(3)** used the integral representation to approximate $S\alpha S$ densities.
- derived fractional moments $\mathbb{E}[|X \mu|^p]$.
- We then briefly showed that these ideas can be extended to the multivariable case.
- Suggested that the complex fractional moments $\mathbb{E}[(-iX_1^{-\gamma_1})(-iX_2^{-\gamma_2})]$ may be useful to help understand dependence structures.